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## ON THE EXISTENCE OF FLOW OF A HEAVY PERFECT FLUID IN A CHANNEL WITH A SLOPING BOTTOM

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Solvability of the problem of a flow of a heavy perfect fluid with free boundary in a channel with a bottom sloping without bounds is proved under the conditions that the Froude number is greater than unity.

Flows in a channel the bottom of which has two horizontal asymptotics, were studied earlier [1-3] under analogous conditions. The problems in which a heavy fluid flows out of a vessel with the rate of flow increasing withour bounds, were studied for the arbitrary [4,5] and for sufficiently large [6] values of the Froude number. However, in all the above problems the free boundary, in contrast to the flow in a channel, passes through a single point at infinity.

Let us consider, in the z = x + iy plane, a steady potential flow of a perfect incompressible heavy fluid with a free boundary, in a channel the bottom of which consists of two rays emerging from the point z = 0 and forming the angles of  $\pi$  and  $-\alpha\pi$  with the positive direction of the x-axis. The velocity vector at  $x \to -$ 

 $\infty$  and the vector of acceleration due to gravity have the corresponding projections  $(v_0, 0)$  and  $(0, -\gamma)$ . Let us introduce the following notation: v and  $\theta$  are the modulus of the velocity vector and the angle between this vector and the x-axis,  $\varphi$  and

 $\psi$  are the velocity potential and stream function,  $w = \varphi + i\psi$ ,  $\psi_0$  is the fluid flow rate,  $\omega = \tau - i\theta = \ln (v_0^{-1}dw / dz)$ ,  $\lambda = \gamma \psi_0 / v_0^{-3}$ .

Mapping the strip  $0 \leqslant \psi \leqslant \psi_0$  in the *w*-variable plane onto the strip  $0 \leqslant \eta \leqslant \pi/2$  in the  $\zeta = \xi + i\eta$  variable plane by means of the function  $w = 2\psi_0\zeta/\pi$  we obtain, from the Bernoulli equation in the usual manner,

$$\exp\left[3\tau\left(\xi+i\frac{\pi}{2}\right)\right] = 1 - \frac{6\lambda}{\pi} \int_{-\infty}^{\xi} \sin\theta\left(\xi+i\frac{\pi}{2}\right) d\xi$$
(1)

Function  $\omega(\zeta)$  satisfies the boundary conditions (1) and the conditions

Using the Will's formulas we obtain

$$u(\xi) = D[g(\xi)] = D_1[g(\xi)] + D_2[g(\xi)]$$

$$D_k[g(\xi)] = -\frac{1}{\pi} \int_{m_k}^{m_k} g(t) \ln \left| th \frac{t-\xi}{2} \right| dt$$

$$k = 1, 2, \quad m_1 = -\infty, \quad m_2 = n_1 = \xi_0, \quad n_2 = \infty$$
(2)

$$\begin{aligned} (u(\xi) &= -\theta(\xi + i\pi/2) - u_0(\xi), \quad u_0(\xi) &= \alpha(\pi - 2 \arctan e^{-\xi}), \\ g(\xi) &= d\tau(\xi + i\pi/2) / d\xi \end{aligned}$$

where  $\xi_0$  is an arbitrary number. The operators  $D_k$  are positive. From (1) and (2) we obtain

$$u(\xi) = D[G(\xi)]$$
(3)  
$$G(\xi) = \frac{2\lambda}{\pi} \sin[u_0(\xi) + u(\xi)] \left[ 1 + \frac{6\lambda}{\pi} \int_{-\infty}^{\xi} \sin[u_0(\xi) + u(\xi)] d\xi \right]^{-1}$$

Equations (3) are equivalent to an operator equation of the type u = T(u), and the solvability of the latter is proved below with help of the Schauder's principle.

Let  $\xi_0 = \ln \sqrt{3}$  and let the parameters  $\alpha$  and  $\lambda$  satisfy the conditions

 $0 < \alpha < 1, \quad 0 \leq \lambda < \min[1,(1-\alpha)\pi - \delta]$ (4)

where  $\delta > 0$  is arbitrarily small (the condition for  $\lambda$  is the same as the inequalities (2.28) and (4.31) in [1]). Let E be a space of functions continuous on  $[-\infty, \infty]$ with the norm  $||u|| = \sup |u(\xi)|$ , and  $H = H(C_1, C_2, \beta)$  be a closed set of E the elements  $u(\xi)$  of which satisfy the conditions

$$0 \leqslant u \ (\xi) \leqslant (1-\alpha)\pi - \delta \quad (|\xi| < \infty) \tag{5}$$

$$\begin{array}{l} u(\xi) \leqslant C_1 e^{-\beta|\xi|} \ (\xi \leqslant \xi_0), \quad u(\xi) \leqslant C_2 / \xi \quad (\xi \geqslant \xi_0) \\ 0 < \beta < 1, \quad C_1 > 0, \quad C_2 > 0 \end{array}$$
(6)

We shall find the estimate of  $F(\xi) = D[G(\xi)] = T(u)$  for  $u \in H$ . We note that  $0 \leq u_0(\xi) \leq \alpha \pi \quad (|\xi| < \infty)$  $u_1(\xi) \leq \alpha \alpha^{-|\xi|} \quad (\xi \leq \xi) \quad u_2(\xi) > \alpha \pi / 2 \quad (\xi > \xi)$ (7)100

$$u_0(\xi) \leqslant 4\alpha e^{-1\varepsilon t} \quad (\xi \leqslant \xi_0), \quad u_0(\xi) \geqslant \alpha \pi / 3 \quad (\xi \geqslant \xi_0) \tag{8}$$

From (5) and (7) it follows that  $0 \leq u_0(\xi) + u(\xi) \leq \pi - \delta$  when  $\xi \in (-\infty, \infty)$  and

$$u_0 + u \geqslant \sin (u_0 + u) \geqslant \pi^{-1} \sin \delta u_0 \geqslant 0 \tag{9}$$

Using (6), (8) and (9) we obtain from (3)

$$0 \leqslant G\left(\xi\right) \leqslant 2\lambda / \pi \quad (|\xi| < \infty) \tag{10}$$

$$0 \leqslant G (\xi) \leqslant 2\lambda / \pi \quad (|\xi| < \infty)$$

$$G (\xi) \leqslant \pi (\alpha \sin \delta \xi)^{-1} \quad (\xi \ge \xi_0)$$

$$G (\xi) \leqslant 8\alpha \pi^{-1} e^{-|\xi|} + 2\lambda \pi^{-1} C_1 e^{-\beta |\xi|} \quad (\xi \leqslant \xi_0)$$
(11)

Applying the operator D to (10) and remembering that D (1) =  $\pi/2$  we obtain  $0 \leq \pi/2$  $F(\xi) \leq \lambda$  and this, together with (4), yields

$$0 \leqslant F(\xi) \leqslant (1-\alpha) \pi - \delta \quad (|\xi| < \infty)$$
<sup>(12)</sup>

Below we shall use the following estimates:

$$D_1\left(e^{-\nu|\xi|}\right) \leqslant \frac{\pi}{2} f\left(\nu\right) e^{-\nu|\xi|} \quad (|\xi| < \infty) \tag{13}$$

$$D_{\mathfrak{s}}(1 / \xi) \leqslant N \Phi(\xi) \quad (|\xi| < \infty) \tag{14}$$

$$\Phi (\xi) = e^{-|\xi|} \quad (\xi \leqslant \xi_0), \quad \Phi (\xi) = 1 / \xi \quad (\xi > \xi_0)$$

where N > 0, f(v) is a continuous increasing function, f(0) = 1 and  $f(1) = \infty$ . The inequality (13) has been proved in [1]. To prove (14), we shall consider the

(10)

function  $r(\zeta) = \ln \ln (e^{\zeta} + 2) - \ln \ln 2$  continuous for  $0 \le \eta \le \pi/2$ ,  $|\xi| < \infty$ . We denote

$$p_{1}(\xi) = \operatorname{Im} r\left(\xi + i\frac{\pi}{2}\right) = \operatorname{arctg} \frac{a}{b}$$

$$p_{2}(\xi) = \frac{d}{d\xi} \operatorname{Re} r\left(\xi + i\frac{\pi}{2}\right) = \frac{2(2a + be^{\xi})e^{\xi}}{(a^{2} + b^{2})(e^{2\xi} + 4)}$$

$$(a = 2 \operatorname{arctg} (e^{\xi} / 2), \quad b = \ln (e^{2\xi} + 4))$$

The following inequalities hold:

 $p_1(\xi) \leqslant M\Phi(\xi) \quad (|\xi| < \infty), \quad p_2(\xi) \ge m/\xi > 0 \quad (\xi \ge \xi_0)$ (15)

Taking into account the fact that  $r(-\infty + i\eta) = 0$ , Im  $r(\xi) = 0$  ( $|\xi| < \infty$ ) we have  $p_1(\xi) = D[p_2(\xi)] \ge D_2[p_2(\xi)]$  and this, together with (15), yields  $M\Phi(\xi) \ge mD_2(1/\xi)$  ( $|\xi| < \infty$ ) (16)

The estimate (14) with N = M / m now follows from (16).

Majorizing  $e^{-|\xi|}$  in (11) with the function  $e^{-\beta|\xi|}$  and applying the estimates (13) and (14) we obtain, from (11),

$$F(\xi) \leqslant \pi N \ (\alpha \sin \delta)^{-1} \Phi(\xi) + \lambda f(\beta) e^{-\beta |\xi|} \ (C_1 + 4\alpha)$$

$$(|\xi| < \infty)$$

$$(17)$$

Using (17), the properties of f(v) and the inequalities  $\lambda < 1$ ,  $0 < \beta < 1$ , we can show that a sufficiently small  $\beta$  and a sufficiently large  $C_1$ , and hence a sufficiently large  $C_2$ , can be chosen such that the following inequalities will hold:

$$F(\xi) \leqslant C_1 e^{-\beta |\xi|} \quad (\xi \leqslant \xi_0), \quad F(\xi) \leqslant \frac{C_2}{\xi} \quad (\xi \geqslant \xi_0) \tag{18}$$

Comparing (5) and (6) with (12) and (18) we conclude, that for the values of  $C_1$ ,  $C_2$  and  $\beta$  chosen the operator T transforms  $H(C_1, C_2, \beta)$  into itself. Using the estimates obtained, we can also show the complete continuity of the contraction of the operator T on H on the norm of the space E (the proof is based on the known fact that the family of functions  $F_n(\xi) = D[G_n(\xi)]$  is equicontinuous on any finite interval provided that the family  $G_n(\xi)$  is uniformly bounded for  $|\xi| < \infty$ ).

From the Schauder principle it follows that when the conditions (4) hold, the equation u = T(u), equivalent to the initial hydrodynamic problem, has at least one solution  $u(\xi) \in H$ . Using the methods of [1, 3] we can extend the theorem of existence proved above to the case of a channel with a curvilinear bottom sloping without bounds.

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